

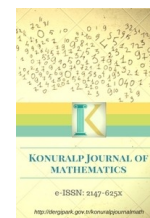
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Riemannian Π –Structure on 5–Dimensional Nilpotent Lie Algebras

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Abstract

The object of our investigations is to classify 5–dimensional nilpotent Lie algebras with two different Riemannian Π –structures. It is shown that the Lie groups corresponding to the Lie algebras \mathfrak{g}_i equipped with two different Riemannian Π –structures are not para-Sasaki-like. Moreover, we investigate whether the considered manifolds admit Ricci-like solitons and whether they are Einstein-like manifolds.

Keywords: Five dimensional nilpotent Lie algebras; Para-Sasaki-like manifold; Ricci-like soliton; Riemannian Π –manifold.

2010 Mathematics Subject Classification: 53C05; 53C15; 53C25.

1. Introduction

The notion of an almost paracontact structure on a smooth odd dimensional manifold was presented in [9, 10]. The geometry of Riemannian manifolds with an almost paracontact structure corresponding to an almost paracomplex structure has been intensively studied in [1, 2, 3, 4, 6]. These manifolds are called briefly Riemannian Π –manifolds. A classification with eleven basic classes of almost paracontact Riemannian manifolds of type (n, n) according to the covariant derivative of the $(1, 1)$ – tensor of the almost paracontact structure was given in [4]. There are 2^{11} classes of Riemannian Π –structures. The investigations of Riemannian Ricci solitons carried out in [7]. Ricci solitons on manifolds such as Riemannian Π – manifolds, Kenmotsu manifolds, paracontact manifolds have been studied in [1, 2, 3, 5, 11].

Non-isomorphic non-abelian nilpotent Lie algebras in five dimensions have six classes [8]. Our aim in this study determine the explicit classes of two different Riemannian Π –structures defined on 5–dimensional nilpotent Lie algebras. Then, we calculate Ricci curvature tensor and scalar curvature tensor. Considering the classification obtained, we see that none of them with given structures are para-Sasaki-like. In addition, we show that the only Lie algebra \mathfrak{g}_1 is an η –Einstein manifold and admits Ricci-like soliton.

The present paper is structured as follows. In Section 2, we reminisce some basic facts and properties of Riemannian Π –manifolds. In Section 3, we classify five dimensional nilpotent Lie algebras with two different Riemannian Π –structure. Finally, we examine some properties of the considered manifolds.

2. Riemannian Π –manifolds

A triple (ϕ, ξ, η) on a $(2n + 1)$ –dimensional smooth manifold M satisfying

$$\phi^2 = Id - \eta \otimes \xi, \quad \eta(\xi) = 1, \quad (2.1)$$

where ϕ is a tensor field of type $(1, 1)$, ξ is a Reeb vector field and η is a 1–form on M , is called an almost paracontact structure on M . In this case, M is called an almost paracontact manifold. In addition, if (M, ϕ, ξ, η) admits a Riemannian metric g with

$$g(\phi x, \phi y) = g(x, y) - \eta(x)\eta(y),$$

for all vector fields x, y , then, (M, ϕ, ξ, η, g) is called Riemannian Π –manifold. These manifolds are sometimes called by different names such as apapR manifolds, almost paracontact almost paracomplex Riemannian manifolds. Moreover, by using above basic identities, the following derived properties are valid:

$$\begin{aligned} g(x, \xi) &= \eta(x), & g(x, \phi y) &= g(\phi x, y), \\ g(\xi, \xi) &= 1, & \eta(\nabla_x \xi) &= 0, \end{aligned} \quad (2.2)$$

where ∇ denotes the Levi-Civita connection of g . The associated metric \tilde{g} of g on (M, ϕ, ξ, η, g) determined by the equality $\tilde{g}(x, y) = g(x, \phi y) + \eta(x)\eta(y)$ is a pseudo-Riemannian metric of signature $(n+1, n)$.

A Riemannian Π -manifold M is said to be a para-Sasaki-like manifold if the following is provided:

$$\begin{aligned}(\nabla_x \phi)y &= -g(x, y)\xi - \eta(y)x + 2\eta(x)\eta(y)\xi, \\ &= -g(\phi x, \phi y)\xi - \eta(y)\phi^2 x.\end{aligned}\quad (2.3)$$

In [6], it is proven that the following identities hold for any para-Sasaki-like manifold (M, g, ϕ, ξ, η) :

$$\begin{aligned}\nabla_x \xi &= \phi x, & (\nabla_x \eta)y &= g(x, \phi y), \\ R(x, y)\xi &= -\eta(y)x + \eta(x)y, & \text{Ric}(x, \xi) &= -2n \eta(x), \\ R(\xi, y)\xi &= \phi^2 y, & \text{Ric}(x, \xi)(\xi, \xi) &= -2n,\end{aligned}\quad (2.4)$$

where R and Ric denote the curvature tensor and the Ricci tensor, respectively.

In [4] the almost paracantact almost paracomplex Riemannian manifolds are classified using the tensor F of type $(0, 3)$ defined by

$$F(x, y, z) = g((\nabla_x \phi)y, z),$$

where ∇ is the Levi-Civita connection of g . Moreover, the following relations are satisfied:

$$\begin{aligned}F(x, y, z) &= F(x, z, y) = -F(x, \phi y, \phi z) + \eta(y)F(x, \xi, z) + \eta(z)F(x, y, \xi), \\ (\nabla_x \eta)y &= g(\nabla_x \xi, y) = -F(x, \phi y, \xi).\end{aligned}\quad (2.5)$$

Eleven basis classes of these manifolds are denoted by $\mathcal{F}_1, \dots, \mathcal{F}_{11}$. The class of \mathcal{F}_0 is defined by the condition $F = 0$, i.e., $\nabla \phi = \nabla \xi = \nabla \eta = 0$.

The Lie 1-forms associated with F are defined by

$$\theta(x) = g^{ij}F(e_i, e_j, x), \quad \theta^*(x) = g^{ij}F(e_i, \phi e_j, x), \quad \omega(x) = F(\xi, \xi, x), \quad (2.6)$$

where g^{ij} 's are the entries of the inverse matrix of g with respect to the basis $\{e_i, \xi\}$ of $T_p M$.

Let \mathbb{F} be the set of all tensors over $T_p M$ satisfying the properties (2.5). \mathbb{F} is the direct sum of eleven subspaces \mathbb{F}_i , which is orthogonal and invariant with respect to the structure group of considered manifolds. If the tensor F belongs to the subspace \mathbb{F}_i , then the manifold is said to be in the class \mathcal{F}_i . It is said that M belongs to the class \mathcal{F}_i if and only if the equality $F = F_i$ is valid. F_i are the components of F in the subspace \mathbb{F}_i and are listed below [4].

$$\begin{aligned}F_1(x, y, z) &= \frac{1}{2n}[g(\phi x, \phi y)\theta(\phi^2 z) + g(\phi x, \phi z)\theta(\phi^2 y) \\ &\quad - g(x, \phi y)\theta(\phi z) - g(x, \phi z)\theta(\phi y)], \\ F_2(x, y, z) &= \frac{1}{4}[2F(\phi^2 x, \phi^2 y, \phi^2 z) + F(\phi^2 y, \phi^2 z, \phi^2 x) + F(\phi^2 z, \phi^2 x, \phi^2 y) \\ &\quad - F(\phi y, \phi z, \phi^2 x) - F(\phi z, \phi y, \phi^2 x)] \\ &\quad - \frac{1}{2n}[g(\phi x, \phi y)\theta(\phi^2 z) + g(\phi x, \phi z)\theta(\phi^2 y) \\ &\quad - g(x, \phi y)\theta(\phi z) - g(x, \phi z)\theta(\phi y)], \\ F_3(x, y, z) &= \frac{1}{4}[2F(\phi^2 x, \phi^2 y, \phi^2 z) - F(\phi^2 y, \phi^2 z, \phi^2 x) - F(\phi^2 z, \phi^2 x, \phi^2 y) \\ &\quad + F(\phi y, \phi z, \phi^2 x) + F(\phi z, \phi y, \phi^2 x)], \\ F_4(x, y, z) &= \frac{\theta(\xi)}{2n}[g(\phi x, \phi y)\eta(z) + g(\phi x, \phi z)\eta(y)], \\ F_5(x, y, z) &= \frac{\theta^*(\xi)}{2n}[g(x, \phi y)\eta(z) + g(x, \phi z)\eta(y)], \\ F_6(x, y, z) &= \frac{1}{4}[[F(\phi^2 x, \phi^2 y, \xi) + F(\phi^2 y, \phi^2 x, \xi) + F(\phi x, \phi y, \xi) + F(\phi y, \phi x, \xi)]\eta(z) \\ &\quad + [F(\phi^2 x, \phi^2 z, \xi) + F(\phi^2 z, \phi^2 x, \xi) + F(\phi x, \phi z, \xi) + F(\phi z, \phi x, \xi)]\eta(y) \\ &\quad - \frac{\theta(\xi)}{2n}[g(\phi x, \phi y)\eta(z) + g(\phi x, \phi z)\eta(y)] \\ &\quad - \frac{\theta^*(\xi)}{2n}[g(x, \phi y)\eta(z) + g(x, \phi z)\eta(y)], \\ F_7(x, y, z) &= \frac{1}{4}[[F(\phi^2 x, \phi^2 y, \xi) - F(\phi^2 y, \phi^2 x, \xi) + F(\phi x, \phi y, \xi) - F(\phi y, \phi x, \xi)]\eta(z) \\ &\quad + [F(\phi^2 x, \phi^2 z, \xi) - F(\phi^2 z, \phi^2 x, \xi) + F(\phi x, \phi z, \xi) - F(\phi z, \phi x, \xi)]\eta(y)], \\ F_8(x, y, z) &= \frac{1}{4}[[F(\phi^2 x, \phi^2 y, \xi) + F(\phi^2 y, \phi^2 x, \xi) - F(\phi x, \phi y, \xi) - F(\phi y, \phi x, \xi)]\eta(z) \\ &\quad + [F(\phi^2 x, \phi^2 z, \xi) + F(\phi^2 z, \phi^2 x, \xi) - F(\phi x, \phi z, \xi) - F(\phi z, \phi x, \xi)]\eta(y)],\end{aligned}$$

$$F_9(x, y, z) = \frac{1}{4} [[F(\phi^2 x, \phi^2 y, \xi) - F(\phi^2 y, \phi^2 x, \xi) - F(\phi x, \phi y, \xi) + F(\phi y, \phi x, \xi)] \eta(z) \\ + [F(\phi^2 x, \phi^2 z, \xi) - F(\phi^2 z, \phi^2 x, \xi) - F(\phi x, \phi z, \xi) + F(\phi z, \phi x, \xi)] \eta(y)],$$

$$F_{10}(x, y, z) = \eta(x) F(\xi, \phi^2 y, \phi^2 z),$$

$$F_{11}(x, y, z) = \eta(x) [\eta(y) \omega(z) + \eta(z) \omega(y)].$$

A Riemannian Π -manifold belongs to a direct sum of two or more basic classes if and only if the fundamental tensor is the sum of the corresponding components F_i, F_j, \dots , namely, $F = F_i + F_j + \dots$.

The Nijenhuis torsion of ϕ is defined by

$$[\phi, \phi](x, y) = [\phi x, \phi y] + \phi^2[x, y] - \phi[\phi x, y] - \phi[x, \phi y]. \quad (2.7)$$

Normality condition of Riemannian Π -structure is equivalent to vanishing the four tensors given by

$$N^{(1)}(x, y) = [\phi, \phi](x, y) - d\eta(x, y)\xi, \\ N^{(2)}(x, y) = (\mathfrak{L}_{\phi x} \eta)(y) - (\mathfrak{L}_{\phi y} \eta)(x), \\ N^{(3)}(x, y) = (\mathfrak{L}_{\xi} \phi)(x), \\ N^{(4)}(x, y) = (\mathfrak{L}_{\xi} \eta)(x),$$

where \mathfrak{L} denotes the Lie derivative operator.

Let us recall from [1] that the Riemannian Π -manifold (M, ϕ, ξ, η, g) is called Einstein-like with constants (a, b, c) if its Ricci tensor Ric satisfies the following formula:

$$\text{Ric} = a g + b \tilde{g} + c \eta \otimes \eta, \quad (2.8)$$

where a, b, c are constants. In particular, if $b = 0$ and $b = c = 0$, then the manifold is called an η -Einstein manifold and an Einstein manifold, respectively. If a, b, c are functions on M , the manifold M is called almost Einstein-like, almost η -Einstein-like or an almost Einstein manifold, respectively.

A Ricci-like soliton with potential vector field ξ and constants (λ, μ, ν) on a Riemannian Π -manifold (M, ϕ, ξ, η, g) is defined by

$$\frac{1}{2} \mathcal{L}_{\xi} g + \text{Ric} + \lambda g + \mu \tilde{g} + \nu \eta \otimes \eta = 0, \quad (2.9)$$

where the Lie derivative \mathcal{L} of g along ξ is expressed by

$$\mathcal{L}_{\xi} g(x, y) = g(\nabla_x \xi, y) + g(x, \nabla_y \xi).$$

An almost paracontact almost paracomplex metric structure (ϕ, ξ, η, g) on a connected Lie group G is said to be left invariant if g is left invariant and the conditions

$$\phi \circ L_a = L_a \circ \phi, \quad L_a(\xi) = \xi$$

are satisfied, where L_a is the left multiplication by $a \in G$ in G .

An almost paracontact almost paracomplex metric structure on G induces an almost paracontact almost paracomplex metric structure on the Lie algebra \mathfrak{g} of G having the structure (ϕ, ξ, η, g) .

In this study, we specify the classes of some almost paracontact almost paracomplex metric structure 5-dimensional nilpotent Lie algebras. The non-isomorphic and non-abelian algebras \mathfrak{g}_i are divided into six classes with the corresponding basis $\{e_1, \dots, e_5\}$ and non-zero brackets in the following [8]:

$$\begin{aligned} \mathfrak{g}_1 &: [e_1, e_2] = e_5, [e_3, e_4] = e_5, \\ \mathfrak{g}_2 &: [e_1, e_2] = e_3, [e_1, e_3] = e_5, [e_2, e_4] = e_5, \\ \mathfrak{g}_3 &: [e_1, e_2] = e_3, [e_1, e_3] = e_4, [e_1, e_4] = e_5, [e_2, e_3] = e_5, \\ \mathfrak{g}_4 &: [e_1, e_2] = e_3, [e_1, e_3] = e_4, [e_1, e_4] = e_5, \\ \mathfrak{g}_5 &: [e_1, e_2] = e_4, [e_1, e_3] = e_5, \\ \mathfrak{g}_6 &: [e_1, e_2] = e_3, [e_1, e_3] = e_4, [e_2, e_3] = e_5. \end{aligned} \quad (2.10)$$

3. A Riemannian Π -structures on 5-dimensional Nilpotent Lie Algebras

Let (ϕ, ξ, η, g) be a left invariant Riemannian Π -structure on a connected Lie group G_i with corresponding Lie algebra \mathfrak{g}_i . We use the same notation for the corresponding Riemannian Π -structure. Now, we investigate the classes of the following Riemannian Π -structure with respect to the basis $\{e_1, \dots, e_5\}$ on each \mathfrak{g}_i .

$$\begin{aligned} \phi(e_1) &= e_3, \phi(e_2) = e_4, \phi(e_3) = e_1, \phi(e_4) = e_2, \phi(e_5) = 0, \\ \xi &= e_5, \eta = e^5, \\ g(e_i, e_i) &= 1, g(e_i, e_j) = 0, i, j \in \{1, \dots, 5\}, i \neq j. \end{aligned} \quad (3.1)$$

3.1. The Lie algebra \mathfrak{g}_1

Theorem 3.1. *The Lie algebra \mathfrak{g}_1 belongs to the class \mathcal{F}_7 according to the structure given in (3.1).*

Proof. By using the non-zero brackets $[e_1, e_2] = e_5$, $[e_3, e_4] = e_5$ and Kozsul's formula, the covariant derivatives of the non-zero basic elements are given by

$$\nabla_{e_1} e_2 = \frac{1}{2} e_5, \nabla_{e_1} e_5 = -\frac{1}{2} e_2, \nabla_{e_2} e_1 = -\frac{1}{2} e_5, \nabla_{e_2} e_5 = \frac{1}{2} e_1,$$

$$\nabla_{e_3} e_4 = \frac{1}{2} e_5, \nabla_{e_3} e_5 = -\frac{1}{2} e_4, \nabla_{e_4} e_3 = -\frac{1}{2} e_5, \nabla_{e_4} e_5 = \frac{1}{2} e_3,$$

$$\nabla_{e_5} e_1 = -\frac{1}{2} e_2, \nabla_{e_5} e_2 = \frac{1}{2} e_1, \nabla_{e_5} e_3 = -\frac{1}{2} e_4, \nabla_{e_5} e_4 = \frac{1}{2} e_3.$$

Theorem 3.2. [4] *Let (M, ϕ, ξ, η, g) be a Riemannian Π -manifold. Then, we have*

- $[\phi, \phi](x, y) = 0$ if and only if (M, ϕ, ξ, η, g) belongs to \mathcal{F}_i ($i = 1, 2, 4, 5, 6, 11$) or to their direct sums;
- $[\phi, \phi](x, y) = -2\{\phi(\nabla_{\phi x} \phi)\phi y + \phi(\nabla_{\phi^2 x} \phi)\phi^2 y\}$ if and only if (M, ϕ, ξ, η, g) belongs to \mathcal{F}_3 ;
- $[\phi, \phi](x, y) = -2(\nabla_x \eta)(y)\xi$ if and only if (M, ϕ, ξ, η, g) belongs to \mathcal{F}_7 ;
- $[\phi, \phi](x, y) = -2\{\eta(x)\nabla_y \xi - \eta(y)\nabla_x \xi - (\nabla_x \eta)(y)\xi\}$ if and only if (M, ϕ, ξ, η, g) belongs to \mathcal{F}_8 ;
- $[\phi, \phi](x, y) = -2\{\eta(x)\nabla_y \xi - \eta(y)\nabla_x \xi\}$ if and only if (M, ϕ, ξ, η, g) belongs to \mathcal{F}_9 ;
- $[\phi, \phi](x, y) = -\eta(x)\phi(\nabla_\xi \phi)y + \eta(y)\phi(\nabla_\xi \phi)x$ if and only if (M, ϕ, ξ, η, g) belongs to \mathcal{F}_{10} .

Setting $x = y = e_i$, ($i = 1, 2, \dots, 5$) in (2.7), we get

$$[\phi e_i, \phi e_i] + \phi^2[e_i, e_i] - \phi[\phi e_i, e_i] - \phi[e_i, \phi e_i] = 0.$$

Moreover, we can calculate

$$(\nabla_{e_i} \eta)e_i = e_i(\eta e_i) - \eta(\nabla_{e_i} e_i) = 0,$$

for every $i = 1, 2, \dots, 5$. In case of $i = 1$, $j = 2$, we obtain $[\phi, \phi](e_1, e_2) = e_5$ and $(\nabla_{e_1} \eta)e_2 = -\frac{1}{2}$. Similarly, in case of $i = 3$, $j = 4$, we get $[\phi, \phi](e_3, e_4) = e_5$ and $(\nabla_{e_3} \eta)e_4 = -\frac{1}{2}$. In other cases, we calculate $[\phi, \phi](e_i, e_j) = 0$ and $(\nabla_{e_i} \eta)(e_j) = 0$ for $i \neq j$. Therefore, the equality given in Theorem 3.2(c) is satisfied for the orthonormal basis $\{e_1, \dots, e_5 = \xi\}$. Hence, we conclude that \mathfrak{g}_1 belongs to the class \mathcal{F}_7 . \square

Now, we consider another structure (ϕ, ξ, η, g) given by

$$\begin{aligned} \phi(e_3) &= e_5, \phi(e_2) = e_4, \phi(e_5) = e_3, \phi(e_4) = e_2, \phi(e_1) = 0, \\ \xi &= e_1, \eta = e^1, \\ g(e_i, e_i) &= 1, g(e_i, e_j) = 0, i, j \in \{1, \dots, 5\}, i \neq j. \end{aligned} \quad (3.2)$$

By using above structure we compute the following non-zero components $F(e_i, e_j, e_k) = F_{ijk}$ of the structure tensor F :

$$\begin{aligned} F_{145} &= F_{154} = F_{213} = F_{231} = F_{325} = \frac{1}{2}, \\ F_{352} &= F_{514} = F_{523} = F_{532} = F_{541} = \frac{1}{2}, \\ F_{123} &= F_{132} = F_{334} = F_{343} = F_{545} = F_{554} = -\frac{1}{2}, \\ F_{433} &= -F_{455} = 1. \end{aligned}$$

Then, we construct the following form of F for any vectors x, y, z :

$$\begin{aligned} F(x, y, z) &= F\left(\sum_i x_i e_i, \sum_j y_j e_j, \sum_k z_k e_k\right) \\ &= \sum_{i,j,k} x_i y_j z_k F(e_i, e_j, e_k) \\ &= -\frac{1}{2} x_1 y_2 z_3 - \frac{1}{2} x_1 y_3 z_2 + \frac{1}{2} x_1 y_4 z_5 + \frac{1}{2} x_1 y_5 z_4 + \frac{1}{2} x_2 y_1 z_3 + \frac{1}{2} x_2 y_3 z_1 \\ &\quad + \frac{1}{2} x_3 y_2 z_5 - \frac{1}{2} x_3 y_3 z_4 - \frac{1}{2} x_3 y_4 z_3 + \frac{1}{2} x_3 y_5 z_2 + x_4 y_3 z_3 - x_4 y_5 z_5 \\ &\quad + \frac{1}{2} x_5 y_1 z_4 + \frac{1}{2} x_5 y_2 z_3 + \frac{1}{2} x_5 y_3 z_2 + \frac{1}{2} x_5 y_4 z_1 - \frac{1}{2} x_5 y_4 z_5 - \frac{1}{2} x_5 y_5 z_4. \end{aligned}$$

The latter equality implies that F is represented in the form

$$F(x, y, z) = F_1(x, y, z) + F_2(x, y, z) + F_3(x, y, z) + F_6(x, y, z) + F_9(x, y, z) + F_{10}(x, y, z),$$

where

$$F_1(x, y, z) = \frac{1}{4}(-x_1y_1z_4 - x_1y_4z_1 + x_3y_2z_5 - x_3y_3z_4 - x_3y_4z_3 + x_3y_5z_2 \\ + 2x_4y_2z_2 + x_5y_2z_3 + x_5y_3z_2 - x_5y_4z_5 - x_5y_5z_4),$$

$$F_2(x, y, z) = \frac{1}{4}(2x_3y_5z_2 + x_4y_3z_3 - 3x_4y_5z_5 + x_5y_2z_3 - x_5y_4z_5 \\ - 2x_5y_5z_4 + x_2y_3z_5 + x_1y_1z_4 + x_1y_4z_1 + x_3y_4z_3 - 2x_4y_4z_2),$$

$$F_3(x, y, z) = \frac{1}{4}(2x_3y_2z_5 - x_3y_5z_2 - x_3y_3z_4 - 2x_3y_4z_3 + 3x_4y_3z_3 \\ - x_4y_5z_5 + x_5y_3z_2 + x_5y_5z_4 - x_2y_3z_5),$$

$$F_6(x, y, z) = \frac{1}{4}(x_2y_3z_1 + x_5y_4z_1 + x_3y_2z_1 + x_4y_5z_1 + x_2y_1z_3 + x_5y_1z_4 + x_3y_1z_2 + x_4y_1z_5),$$

$$F_9(x, y, z) = \frac{1}{4}(x_2y_3z_1 + x_5y_4z_1 - x_3y_2z_1 - x_4y_5z_1 + x_2y_1z_3 + x_5y_1z_4 - x_3y_1z_2 - x_4y_1z_5),$$

$$F_{10}(x, y, z) = -\frac{1}{2}x_1y_2z_3 - \frac{1}{2}x_1y_3z_2 + \frac{1}{2}x_1y_4z_5 + \frac{1}{2}x_1y_5z_4.$$

Therefore, \mathfrak{g}_1 with the structure (3.2) is in the class $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_6 \oplus \mathcal{F}_9 \oplus \mathcal{F}_{10}$.

The Ricci tensor Ric and the scalar curvature scal according to the basis $\{e_1, \dots, e_4, e_5 = \xi\}$ are presented by

$$\text{Ric}(x, y) = \sum_{i=1}^5 g(R(e_i, x)y, e_i) \text{ and } \text{scal} = \sum_{i=1}^5 \text{Ric}(e_i, e_i), \quad (3.3)$$

respectively. The non-zero components of Ricci tensor Ric corresponding to the Lie algebra \mathfrak{g}_1 are calculated according to the basis $\{e_1, \dots, e_4, e_5 = \xi\}$ as follows:

$$\begin{aligned} Ric_{11} &= -\frac{1}{2}, & Ric_{22} &= -\frac{1}{2}, \\ Ric_{33} &= -\frac{1}{2}, & Ric_{44} &= -\frac{1}{2}, \\ Ric_{55} &= 1, \end{aligned}$$

where $Ric_{ij} = Ric(e_i, e_j)$ for $i, j \in \{1, 2, \dots, 5\}$. The scalar curvature scal of \mathfrak{g}_1 is evaluated by $\text{scal} = -1$. $(G_1, \phi, \xi, \eta, g)$ is a η -Einstein manifold with constants $(a, b, c) = (-\frac{1}{2}, 0, \frac{3}{2})$.

The nonzero components of $\mathcal{L}_\xi g$ for the structure (3.2) are the following:

$$(\mathcal{L}_\xi g)_{25} = (\mathcal{L}_\xi g)_{52} = -1.$$

$(G_1, \phi, \xi, \eta, g)$ is neither Einstein-like nor Ricci-like soliton for the structure (3.2).

3.2. The Lie algebra \mathfrak{g}_2

Theorem 3.3. *The Lie algebra \mathfrak{g}_2 belongs to the class $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$ with regard to the structure given in (3.1).*

Proof. By using the relations $[e_1, e_2] = e_3, [e_1, e_3] = e_5, [e_2, e_4] = e_5$ and Kozsul's formula we get

$$\nabla_{e_1}e_2 = \frac{1}{2}e_3, \nabla_{e_1}e_3 = -\frac{1}{2}e_2 + \frac{1}{2}e_5, \nabla_{e_1}e_5 = -\frac{1}{2}e_3, \nabla_{e_2}e_1 = -\frac{1}{2}e_3,$$

$$\nabla_{e_2}e_3 = \frac{1}{2}e_1, \nabla_{e_2}e_4 = \frac{1}{2}e_5, \nabla_{e_2}e_5 = -\frac{1}{2}e_4, \nabla_{e_3}e_1 = -\frac{1}{2}e_2 - \frac{1}{2}e_5,$$

$$\nabla_{e_3}e_2 = \frac{1}{2}e_1, \nabla_{e_3}e_5 = \frac{1}{2}e_1, \nabla_{e_4}e_2 = -\frac{1}{2}e_5, \nabla_{e_4}e_5 = \frac{1}{2}e_2,$$

$$\nabla_{e_5}e_1 = -\frac{1}{2}e_3, \nabla_{e_5}e_2 = -\frac{1}{2}e_4, \nabla_{e_5}e_3 = \frac{1}{2}e_1, \nabla_{e_5}e_4 = \frac{1}{2}e_2.$$

We evaluate the projections and determine the class of the structure. The nonzero structure constants F_{ijk} are given in the following:

$$\begin{aligned} F_{115} &= F_{134} = F_{143} = F_{151} = \frac{1}{2}, \\ F_{225} &= F_{252} = F_{314} = F_{341} = \frac{1}{2}, \\ F_{112} &= F_{121} = F_{323} = F_{332} = -\frac{1}{2}, \\ F_{335} &= F_{353} = F_{445} = F_{454} = -\frac{1}{2}, \\ F_{211} &= F_{511} = F_{522} = 1, \\ F_{233} &= F_{533} = F_{544} = -1. \end{aligned}$$

For any x, y, z , by using above relations, the tensor F can be calculated in the following way:

$$\begin{aligned}
 F(x, y, z) &= F\left(\sum_i x_i e_i, \sum_j y_j e_j, \sum_k z_k e_k\right) \\
 &= \sum_{i,j,k} x_i y_j z_k F(e_i, e_j, e_k) \\
 &= -\frac{1}{2}x_1 y_1 z_2 - \frac{1}{2}x_1 y_2 z_1 + \frac{1}{2}x_1 y_3 z_4 + \frac{1}{2}x_1 y_4 z_3 + \frac{1}{2}x_3 y_4 z_1 + \frac{1}{2}x_3 y_1 z_4 \\
 &\quad + \frac{1}{2}x_1 y_1 z_5 + \frac{1}{2}x_2 y_2 z_5 + x_2 y_1 z_1 - x_2 y_3 z_3 - \frac{1}{2}x_3 y_3 z_5 - \frac{1}{2}x_4 y_4 z_5 \\
 &\quad + \frac{1}{2}x_1 y_5 z_1 + \frac{1}{2}x_2 y_5 z_2 - \frac{1}{2}x_3 y_2 z_3 - \frac{1}{2}x_3 y_5 z_3 - \frac{1}{2}x_4 y_5 z_4 - \frac{1}{2}x_3 y_3 z_2 \\
 &\quad + x_5 y_1 z_1 + x_5 y_2 z_2 - x_5 y_3 z_3 - x_5 y_4 z_4.
 \end{aligned}$$

Since

$$\begin{aligned}
 F_1(x, y, z) &= \frac{1}{4}(-x_1 y_1 z_2 - 2x_2 y_2 z_2 - x_3 y_3 z_2 - x_5 y_5 z_2 - x_1 y_2 z_1 - x_3 y_2 z_3 \\
 &\quad - x_5 y_2 z_5 + x_1 y_4 z_3 + 2x_2 y_4 z_4 + x_3 y_4 z_1 + x_1 y_3 z_4 + x_3 y_1 z_4), \\
 F_2(x, y, z) &= \frac{1}{4}(2x_2 y_1 z_1 + x_2 y_2 z_2 - 2x_2 y_3 z_3 - 2x_2 y_4 z_4 + 2x_3 y_1 z_4 \\
 &\quad + 2x_3 y_4 z_1 - 2x_3 y_2 z_3 - 2x_3 y_3 z_2 + x_5 y_2 z_5 + x_5 y_5 z_2), \\
 F_3(x, y, z) &= \frac{1}{4}(-x_1 y_1 z_2 - x_1 y_2 z_1 + x_1 y_3 z_4 + x_1 y_4 z_3 + 2x_2 y_1 z_1 \\
 &\quad - 2x_2 y_3 z_3 - x_3 y_1 z_4 + x_3 y_2 z_3 + x_3 y_3 z_2 - x_3 y_4 z_1), \\
 F_8(x, y, z) &= \frac{1}{2}x_1 y_1 z_5 + \frac{1}{2}x_2 y_2 z_5 - \frac{1}{2}x_3 y_3 z_5 - \frac{1}{2}x_4 y_4 z_5 + \frac{1}{2}x_1 y_5 z_1 \\
 &\quad + \frac{1}{2}x_2 y_5 z_2 - \frac{1}{2}x_3 y_5 z_3 - \frac{1}{2}x_4 y_5 z_4, \\
 F_{10}(x, y, z) &= x_5 y_1 z_1 + x_5 y_2 z_2 - x_5 y_3 z_3 - x_5 y_4 z_4,
 \end{aligned}$$

the tensor F can be written as $F = F_1 + F_2 + F_3 + F_8 + F_{10}$. The only nonzero projections are $\mathcal{F}_1, \mathcal{F}_2, \mathcal{F}_3, \mathcal{F}_8, \mathcal{F}_{10}$. Therefore, the Lie algebra \mathfrak{g}_2 is in the class $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$. \square

For the structure (3.2), the non-zero components F_{ijk} can be found as

$$\begin{aligned}
 F_{134} &= F_{215} = F_{225} = F_{251} = F_{252} = F_{313} = F_{314} = \frac{1}{2}, \\
 F_{331} &= F_{341} = F_{423} = F_{432} = F_{143} = F_{515} = \frac{1}{2}, \\
 F_{125} &= F_{152} = F_{234} = F_{243} = F_{445} = F_{454} = -\frac{1}{2}, \\
 F_{155} &= F_{522} = -F_{133} = -F_{544} = 1.
 \end{aligned}$$

The only nonzero projections of the tensor F are calculated by

$$\begin{aligned}
 F_2(x, y, z) &= \frac{1}{2}x_2 y_2 z_5 - \frac{1}{2}x_2 y_3 z_4 - \frac{1}{2}x_2 y_4 z_3 + \frac{1}{2}x_2 y_5 z_2 + \frac{1}{2}x_4 y_2 z_3 \\
 &\quad + \frac{1}{2}x_4 y_3 z_2 - \frac{1}{2}x_4 y_4 z_5 - \frac{1}{2}x_4 y_5 z_4 + x_5 y_2 z_2 - x_5 y_4 z_5, \\
 F_4(x, y, z) &= \frac{1}{4}(x_2 y_2 z_1 + x_3 y_3 z_1 + x_4 y_4 z_1 + x_5 y_5 z_1 + x_2 y_1 z_2 + x_3 y_1 z_3 + x_4 y_1 z_4 + x_5 y_1 z_5), \\
 F_6(x, y, z) &= \frac{1}{4}(x_5 y_5 z_1 + x_2 y_5 z_1 + x_3 y_3 z_1 + x_3 y_4 z_1 + x_5 y_2 z_1 + x_4 y_3 z_1 + x_5 y_1 z_5 + x_2 y_1 z_5 \\
 &\quad + x_3 y_1 z_3 + x_3 y_1 z_4 - x_5 y_1 z_2 + x_4 y_1 z_3 - x_2 y_2 z_1 - x_4 y_4 z_1 - x_2 y_1 z_2 - x_4 y_1 z_4),
 \end{aligned}$$

$$F_9(x, y, z) = \frac{1}{4}(x_2y_5z_1 + x_3y_4z_1 - x_5y_2z_1 - x_4y_3z_1 + x_2y_1z_5 + x_3y_1z_4 - x_5y_1z_2 - x_4y_1z_3),$$

$$F_{10}(x, y, z) = -\frac{1}{2}x_1y_2z_5 - x_1y_3z_3 + \frac{1}{2}x_1y_3z_4 + \frac{1}{2}x_1y_4z_3 - \frac{1}{2}x_1y_5z_2 + x_1y_5z_5.$$

Hence, in similar way, it can be easily seen that the structure (3.2) on \mathfrak{g}_2 is of type $\mathcal{F}_2 \oplus \mathcal{F}_4 \oplus \mathcal{F}_6 \oplus \mathcal{F}_9 \oplus \mathcal{F}_{10}$.

The nonzero components $Ric_{ij} = Ric(e_i, e_j)$ of Ricci curvature tensor are given by

$$\begin{aligned} Ric_{11} &= -1, \quad Ric_{22} = -1, \\ Ric_{33} &= 0, \quad Ric_{44} = -\frac{1}{2}, \\ Ric_{55} &= 1. \end{aligned}$$

With the aid of the above relations, we can compute the scalar curvature scal as follows:

$$\begin{aligned} \text{scal} &= \sum_{i=1}^5 Ric_{ii} \\ &= Ric_{11} + Ric_{22} + Ric_{33} + Ric_{44} + Ric_{55} \\ &= -1 - 1 + 0 - \frac{1}{2} + 1 \\ &= -\frac{3}{2} \end{aligned}$$

By direct computation it can be easily shown that the Lie algebra \mathfrak{g}_2 is not Einstein - like manifold. Moreover, the nonzero components of the Lie derivative $\mathcal{L}_\xi g$ for the structure (3.2) are as follows:

$$(\mathcal{L}_\xi g)_{23} = (\mathcal{L}_\xi g)_{35} = (\mathcal{L}_\xi g)_{32} = (\mathcal{L}_\xi g)_{53} = -1.$$

Hence, \mathfrak{g}_2 is not Ricci-like soliton for both structures.

3.3. The Lie algebra \mathfrak{g}_3

Theorem 3.4. *The Lie algebra \mathfrak{g}_3 belongs to the class $\mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$ with respect to the structure given in (3.1).*

Proof. With the aid of the relations given in \mathfrak{g}_3 , the basic components of the Levi-Civita connection ∇ can be found as

$$\begin{aligned} \nabla_{e_1}e_2 &= \frac{1}{2}e_3, \quad \nabla_{e_1}e_3 = -\frac{1}{2}e_2 + \frac{1}{2}e_4, \quad \nabla_{e_1}e_4 = -\frac{1}{2}e_3 + \frac{1}{2}e_5, \quad \nabla_{e_1}e_5 = -\frac{1}{2}e_4, \\ \nabla_{e_2}e_1 &= -\frac{1}{2}e_3, \quad \nabla_{e_2}e_3 = \frac{1}{2}e_1 + \frac{1}{2}e_5, \quad \nabla_{e_2}e_5 = -\frac{1}{2}e_3, \quad \nabla_{e_3}e_1 = -\frac{1}{2}e_2 - \frac{1}{2}e_4, \\ \nabla_{e_3}e_2 &= \frac{1}{2}e_1 - \frac{1}{2}e_5, \quad \nabla_{e_3}e_4 = \frac{1}{2}e_1, \quad \nabla_{e_3}e_5 = \frac{1}{2}e_2, \quad \nabla_{e_4}e_1 = -\frac{1}{2}e_3 - \frac{1}{2}e_5, \quad \nabla_{e_4}e_3 = \frac{1}{2}e_1, \\ \nabla_{e_4}e_5 &= \frac{1}{2}e_1, \quad \nabla_{e_5}e_1 = -\frac{1}{2}e_4, \quad \nabla_{e_5}e_2 = -\frac{1}{2}e_3, \quad \nabla_{e_5}e_3 = \frac{1}{2}e_2, \quad \nabla_{e_5}e_4 = \frac{1}{2}e_1. \end{aligned}$$

By direct computation we get nonzero basic components F_{ijk} of the tensor F as follows:

$$\begin{aligned} F_{114} &= F_{125} = F_{134} = F_{141} = F_{143} = F_{251} = \frac{1}{2}, \\ F_{152} &= F_{215} = F_{312} = F_{314} = F_{321} = F_{341} = \frac{1}{2}, \\ F_{112} &= F_{121} = F_{123} = F_{132} = F_{323} = F_{332} = -\frac{1}{2}, \\ F_{334} &= F_{343} = F_{345} = F_{354} = F_{435} = F_{453} = -\frac{1}{2}, \\ F_{211} &= F_{411} = F_{512} = F_{521} = 1, \\ F_{233} &= F_{433} = F_{534} = F_{543} = -1. \end{aligned}$$

The nonzero projections F_i are in the following:

$$\begin{aligned} F_2(x, y, z) &= \frac{1}{4}(-x_1y_1z_2 + 2x_1y_1z_4 - x_1y_2z_1 - 2x_1y_2z_3 - 2x_1y_3z_2 \\ &\quad + 2x_1y_4z_1 + x_1y_4z_3 + 2x_2y_1z_1 - 2x_2y_3z_3 + 2x_3y_1z_2 + 3x_3y_1z_4 + 2x_3y_2z_1 \\ &\quad - 3x_3y_2z_3 - 3x_3y_3z_2 - 2x_3y_3z_4 + 3x_3y_4z_1 - 2x_3y_4z_3 + 4x_4y_1z_1 - 4x_4y_3z_3), \end{aligned}$$

$$\begin{aligned}
F_3(x, y, z) &= \frac{1}{4}(-x_1y_1z_2 - x_1y_2z_1 + 2x_1y_3z_4 + x_1y_4z_3 + 2x_2y_1z_1 \\
&\quad - 2x_2y_3z_3 - x_3y_1z_4 + x_3y_2z_3 - x_3y_1z_4 + x_3y_2z_3 + x_3y_3z_2 - x_3y_4z_1), \\
F_8(x, y, z) &= \frac{1}{2}x_1y_2z_5 + \frac{1}{2}x_1y_5z_2 + \frac{1}{2}x_2y_1z_5 + \frac{1}{2}x_2y_5z_1 - \frac{1}{2}x_3y_4z_5 \\
&\quad - \frac{1}{2}x_3y_5z_4 - \frac{1}{2}x_4y_3z_5 - \frac{1}{2}x_4y_5z_3, \\
F_{10}(x, y, z) &= x_5y_1z_2 + x_5y_2z_1 - x_5y_3z_4 - x_5y_4z_3.
\end{aligned}$$

Then, the tensor F can be written as

$$F(x, y, z) = F_2(x, y, z) + F_3(x, y, z) + F_8(x, y, z) + F_{10}(x, y, z).$$

The class of \mathfrak{g}_3 according to the structure (3.1) is $\mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$. □

By the structure (3.2), the nonzero basic components F_{ijk} are calculated by

$$\begin{aligned}
F_{145} &= F_{154} = F_{215} = F_{251} = F_{312} = F_{314} = \frac{1}{2}, \\
F_{321} &= F_{323} = F_{332} = F_{341} = F_{413} = F_{415} = \frac{1}{2}, \\
F_{431} &= F_{451} = F_{512} = F_{521} = F_{525} = F_{552} = \frac{1}{2}, \\
F_{123} &= F_{132} = F_{345} = F_{354} = F_{534} = F_{543} = -\frac{1}{2}, \\
F_{255} &= -F_{233} = 1.
\end{aligned}$$

By using above basic components of the tensor F , we obtain the following projections:

$$\begin{aligned}
F_1(x, y, z) &= \frac{1}{4}(2x_2y_2z_2 + x_3y_3z_2 + x_5y_5z_2 + x_3y_2z_3 + x_5y_2z_5 \\
&\quad - 2x_2y_4z_4 - x_3y_5z_4 - x_5y_3z_4 - x_5y_4z_3 - x_3y_4z_5), \\
F_2(x, y, z) &= \frac{1}{2}(-x_2y_3z_3 + x_2y_5z_5 - x_5y_3z_4 - x_5y_4z_3 + x_5y_2z_5 + x_5y_5z_2 - x_2y_2z_2 - x_2y_4z_4) \\
F_3(x, y, z) &= \frac{1}{4}(-2x_2y_3z_3 + 2x_2y_5z_5 + x_3y_2z_3 + x_3y_3z_2 - x_3y_4z_5 \\
&\quad - x_3y_5z_4 + x_5y_3z_4 + x_5y_4z_3 - x_5y_2z_5 - x_5y_5z_2) \\
F_6(x, y, z) &= \frac{1}{4}(2x_2y_5z_1 + x_3y_2z_1 + 2x_3y_4z_1 + 2x_4y_3z_1 + x_4y_5z_1 + 2x_5y_2z_1 + x_2y_3z_1 + x_5y_4z_1 \\
&\quad + 2x_2y_1z_5 + x_3y_1z_2 + 2x_3y_1z_4 + 2x_4y_1z_3 + x_4y_1z_5 + 2x_5y_1z_2 + x_2y_1z_3 + x_5y_1z_4), \\
F_9(x, y, z) &= \frac{1}{4}(x_3y_2z_1 + x_4y_5z_1 - x_2y_3z_1 - x_5y_4z_1 + x_3y_1z_2 + x_4y_1z_5 - x_2y_1z_3 - x_5y_1z_4), \\
F_{10}(x, y, z) &= -\frac{1}{2}x_1y_2z_3 - \frac{1}{2}x_1y_3z_2 + \frac{1}{2}x_1y_4z_5 + \frac{1}{2}x_1y_5z_4.
\end{aligned}$$

Namely, \mathfrak{g}_3 belongs to $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_6 \oplus \mathcal{F}_9 \oplus \mathcal{F}_{10}$.

The nonzero components of Ricci curvature tensor for \mathfrak{g}_3 are given below.

$$\begin{aligned}
Ric_{11} &= -\frac{3}{2}, \quad Ric_{22} = -1, \\
Ric_{33} &= -\frac{1}{2}, \quad Ric_{44} = 0, \\
Ric_{55} &= 1.
\end{aligned}$$

Using above equations, we compute $scal = -2$. The nonzero components of $\mathcal{L}_\xi g$ for the structure (3.2) are the following:

$$(\mathcal{L}_\xi g)_{23} = (\mathcal{L}_\xi g)_{32} = (\mathcal{L}_\xi g)_{34} = (\mathcal{L}_\xi g)_{43} = (\mathcal{L}_\xi g)_{45} = (\mathcal{L}_\xi g)_{54} = -1.$$

By direct calculation, it is easily checked that \mathfrak{g}_3 is not Einstein - like and Ricci-like soliton for given two structure.

3.4. The Lie algebra \mathfrak{g}_4

Theorem 3.5. The Lie algebras \mathfrak{g}_4 belongs to the class $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_7 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$ according to the structure given in (3.1).

Proof. Similarly, by using the relations $[e_1, e_2] = e_3, [e_1, e_3] = e_4, [e_1, e_4] = e_5$ and Kozsul's formula, the basic components of ∇ are calculated by

$$\nabla_{e_1} e_2 = \frac{1}{2} e_3, \nabla_{e_1} e_3 = -\frac{1}{2} e_2 + \frac{1}{2} e_4, \nabla_{e_1} e_4 = -\frac{1}{2} e_3 + \frac{1}{2} e_5, \nabla_{e_1} e_5 = -\frac{1}{2} e_4,$$

$$\nabla_{e_2} e_1 = -\frac{1}{2} e_3, \nabla_{e_2} e_3 = \frac{1}{2} e_1, \nabla_{e_3} e_1 = -\frac{1}{2} e_2 - \frac{1}{2} e_4,$$

$$\nabla_{e_3} e_2 = \frac{1}{2} e_1, \nabla_{e_3} e_4 = \frac{1}{2} e_1, \nabla_{e_4} e_1 = -\frac{1}{2} e_3 - \frac{1}{2} e_5, \nabla_{e_4} e_3 = \frac{1}{2} e_1,$$

$$\nabla_{e_4} e_5 = \frac{1}{2} e_1, \nabla_{e_5} e_1 = -\frac{1}{2} e_4, \nabla_{e_5} e_4 = \frac{1}{2} e_1.$$

The basic components F_{ijk} are given by

$$\begin{aligned} F_{114} &= F_{125} = F_{134} = F_{141} = F_{143} = F_{512} = \frac{1}{2}, \\ F_{152} &= F_{521} = F_{312} = F_{314} = F_{321} = F_{341} = \frac{1}{2}, \\ F_{112} &= F_{121} = F_{123} = F_{132} = F_{323} = F_{332} = -\frac{1}{2}, \\ F_{334} &= F_{343} = F_{435} = F_{453} = F_{534} = F_{543} = -\frac{1}{2}, \\ F_{211} &= F_{411} = -F_{233} = -F_{433} = 1. \end{aligned}$$

Since the projections F_i are too long, they are not written explicitly. It can be seen that the class of \mathfrak{g}_4 is in $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_7 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$. \square

Using the structure given in (3.2), the nonzero structure constants F_{ijk} are given below.

$$\begin{aligned} F_{145} &= F_{154} = F_{215} = F_{251} = F_{312} = F_{314} = \frac{1}{2}, \\ F_{321} &= F_{341} = F_{413} = F_{415} = \frac{1}{2}, \\ F_{431} &= F_{451} = F_{512} = F_{521} = \frac{1}{2}, \\ F_{123} &= F_{132} = -\frac{1}{2}. \end{aligned}$$

Using above relations, we have

$$\begin{aligned} F_6(x, y, z) &= \frac{1}{4} (2x_2y_5z_1 + x_3y_2z_1 + 2x_3y_4z_1 + 2x_4y_3z_1 + x_4y_5z_1 + 2x_5y_2z_1 + x_2y_3z_1 + x_5y_4z_1 \\ &\quad + 2x_2y_1z_5 + x_3y_1z_2 + 2x_3y_1z_4 + 2x_4y_1z_3 + x_4y_1z_5 + 2x_5y_1z_2 + x_2y_1z_3 + x_5y_1z_4), \\ F_9(x, y, z) &= \frac{1}{4} (x_3y_2z_1 + x_4y_5z_1 - x_2y_3z_1 - x_5y_4z_1 + x_3y_1z_2 + x_4y_1z_5 - x_2y_1z_3 - x_5y_1z_4), \\ F_{10}(x, y, z) &= -\frac{1}{2} x_1y_2z_3 - \frac{1}{2} x_1y_3z_2 + \frac{1}{2} x_1y_4z_5 + \frac{1}{2} x_1y_5z_4. \end{aligned}$$

Therefore, we acquire that \mathfrak{g}_4 is in the class $\mathcal{F}_6 \oplus \mathcal{F}_9 \oplus \mathcal{F}_{10}$.

The nonzero components Ric_{ij} of Ricci curvature tensor are determined by the following equations:

$$\begin{aligned} Ric_{11} &= -\frac{3}{2}, \quad Ric_{22} = -\frac{1}{2}, \\ Ric_{55} &= \frac{1}{2}. \end{aligned} \tag{3.4}$$

Taking into account (3.4), we obtain $scal = -\frac{3}{2}$. The nonzero components of $\mathcal{L}_\xi g$ for the structure (3.2) are the following:

$$(\mathcal{L}_\xi g)_{23} = (\mathcal{L}_\xi g)_{32} = (\mathcal{L}_\xi g)_{34} = (\mathcal{L}_\xi g)_{43} = (\mathcal{L}_\xi g)_{45} = (\mathcal{L}_\xi g)_{54} = -1.$$

It can be easily checked that \mathfrak{g}_4 is neither η -Einstein-like nor Ricci-like soliton for given two structures.

3.5. The Lie algebra \mathfrak{g}_5

Theorem 3.6. The class of the Lie algebra \mathfrak{g}_5 is $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$ considering the structure given in (3.1).

Proof. The basic terms of ∇ are computed as follows:

$$\nabla_{e_1}e_2 = \frac{1}{2}e_4, \nabla_{e_1}e_3 = \frac{1}{2}e_5, \nabla_{e_1}e_4 = -\frac{1}{2}e_2, \nabla_{e_1}e_5 = -\frac{1}{2}e_3, \nabla_{e_2}e_1 = -\frac{1}{2}e_4,$$

$$\nabla_{e_2}e_4 = \frac{1}{2}e_1, \nabla_{e_3}e_1 = -\frac{1}{2}e_5, \nabla_{e_3}e_5 = \frac{1}{2}e_1,$$

$$\nabla_{e_4}e_1 = -\frac{1}{2}e_2, \nabla_{e_4}e_2 = \frac{1}{2}e_1, \nabla_{e_5}e_1 = -\frac{1}{2}e_3, \nabla_{e_5}e_3 = \frac{1}{2}e_1.$$

The nonzero projections F_i are given as follows:

$$\begin{aligned} F_1(x, y, z) = & \frac{1}{4}(2x_1y_1z_1 + x_2y_2z_1 + x_3y_3z_1 + x_4y_4z_1 + x_5y_5z_1 + x_2y_1z_2 \\ & + x_3y_1z_3 + x_4y_1z_4 + x_5y_1z_5 - x_2y_4z_3 - x_3y_1z_3 - x_4y_2z_3) \\ & - 2x_1y_3z_3 - x_2y_3z_4 - x_3y_3z_1 - x_4y_3z_2), \end{aligned}$$

$$\begin{aligned} F_2(x, y, z) = & \frac{1}{4}(-2x_1y_2z_2 + 2x_1y_4z_4 - 2x_4y_3z_2 + 2x_4y_4z_1 + 2x_4y_1z_4 - 2x_4y_2z_3 \\ & - 2x_1y_1z_1 - x_5y_5z_1 - x_5y_1z_5 + 2x_1y_3z_3), \end{aligned}$$

$$\begin{aligned} F_3(x, y, z) = & \frac{1}{4}(-2x_1y_2z_2 + 2x_1y_4z_4 + x_2y_1z_2 + x_2y_2z_1 - x_2y_3z_4 \\ & - x_2y_4z_3 + x_4y_3z_2 - x_4y_4z_1 - x_4y_1z_4 + x_4y_2z_3), \end{aligned}$$

$$F_8(x, y, z) = \frac{1}{2}x_1y_1z_5 - \frac{1}{2}x_3y_3z_5 - \frac{1}{2}x_3y_5z_3 + \frac{1}{2}x_1y_5z_1,$$

$$F_{10}(x, y, z) = x_5y_1z_1 - x_5y_3z_3.$$

The basic components of F are calculated by

$$\begin{aligned} F_{115} = F_{151} = F_{212} = F_{221} = F_{414} = F_{441} &= \frac{1}{2}, \\ F_{234} = F_{243} = F_{335} = F_{353} = F_{423} = F_{432} &= -\frac{1}{2}, \\ F_{144} = F_{511} = -F_{122} = -F_{533} &= 1. \end{aligned}$$

If the tensor F is written explicitly for any vectors x, y, z , then we obtain that \mathfrak{g}_5 is in the class $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$. □

Moreover, using the structure (3.2), we have

$$\begin{aligned} F_{212} = F_{221} = F_{313} = F_{331} &= \frac{1}{2}, \\ F_{414} = F_{441} = F_{515} = F_{551} &= \frac{1}{2}, \\ F_{144} = F_{155} = -F_{122} = -F_{133} &= 1. \end{aligned}$$

By direct computation, we get the nonzero projections F_6 and F_{10} in the following way:

$$\begin{aligned} F_6(x, y, z) = & \frac{1}{2}x_2y_1z_2 + \frac{1}{2}x_2y_2z_1 + \frac{1}{2}x_3y_1z_3 + \frac{1}{2}x_3y_3z_1 \\ & + \frac{1}{2}x_4y_1z_4 + \frac{1}{2}x_4y_4z_1 + \frac{1}{2}x_5y_1z_5 + \frac{1}{2}x_5y_5z_1, \end{aligned}$$

$$F_{10}(x, y, z) = -x_1y_2z_2 - x_1y_3z_3 + x_1y_4z_4 + x_1y_5z_5.$$

Hence, we obtain that \mathfrak{g}_5 is in $\mathcal{F}_6 \oplus \mathcal{F}_{10}$.

The non-zero components Ric_{ij} for \mathfrak{g}_5 are

$$\begin{aligned} Ric_{11} = -1, Ric_{22} &= -\frac{1}{2}, \\ Ric_{33} = -\frac{1}{2}, Ric_{44} &= \frac{1}{2}, \\ Ric_{55} &= \frac{1}{2}. \end{aligned}$$

Using above equations, the scalar curvature is -1 . The nonzero components of $\mathcal{L}_\xi g$ for the structure (3.2) are the following:

$$(\mathcal{L}_\xi g)_{24} = (\mathcal{L}_\xi g)_{42} = (\mathcal{L}_\xi g)_{35} = (\mathcal{L}_\xi g)_{53} = -1.$$

It can be easily checked that \mathfrak{g}_5 is neither Einstein-like nor Ricci-like soliton for structures (3.1) and (3.2).

3.6. The Lie algebra \mathfrak{g}_6

Theorem 3.7. *The Lie algebra \mathfrak{g}_6 belongs to the class $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_7 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$ according to the structure given in (3.1).*

Proof. Similarly, the basic components of ∇ are computed as follows:

$$\begin{aligned}\nabla_{e_1}e_2 &= \frac{1}{2}e_3, \nabla_{e_1}e_3 = -\frac{1}{2}e_2 + \frac{1}{2}e_4, \nabla_{e_1}e_4 = -\frac{1}{2}e_3, \\ \nabla_{e_2}e_1 &= -\frac{1}{2}e_3, \nabla_{e_2}e_3 = \frac{1}{2}e_1 + \frac{1}{2}e_5, \nabla_{e_2}e_5 = -\frac{1}{2}e_3, \nabla_{e_3}e_1 = -\frac{1}{2}e_2 - \frac{1}{2}e_4, \\ \nabla_{e_3}e_2 &= \frac{1}{2}e_1 - \frac{1}{2}e_5, \nabla_{e_3}e_4 = \frac{1}{2}e_1, \nabla_{e_3}e_5 = \frac{1}{2}e_2, \nabla_{e_4}e_1 = -\frac{1}{2}e_3, \nabla_{e_4}e_3 = \frac{1}{2}e_1, \\ \nabla_{e_5}e_2 &= -\frac{1}{2}e_3, \nabla_{e_5}e_3 = \frac{1}{2}e_2.\end{aligned}$$

The nonzero components of the structure tensor F are as follows:

$$\begin{aligned}F_{114} &= F_{134} = F_{141} = F_{143} = F_{215} = F_{251} = \frac{1}{2}, \\ F_{312} &= F_{314} = F_{321} = F_{341} = F_{512} = F_{521} = \frac{1}{2}, \\ F_{112} &= F_{121} = F_{123} = F_{132} = F_{323} = F_{332} = -\frac{1}{2}, \\ F_{334} &= F_{343} = F_{345} = F_{354} = F_{534} = F_{543} = -\frac{1}{2}, \\ F_{211} &= F_{411} = -F_{233} = -F_{433} = 1.\end{aligned}$$

We omit the nonzero projections F_i since they are very long. Hence, it is not hard to verify that \mathfrak{g}_6 is in the class $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_7 \oplus \mathcal{F}_8 \oplus \mathcal{F}_{10}$. \square

For the structure (3.2) on \mathfrak{g}_6 , the nonzero components F_{ijk} are in the following:

$$\begin{aligned}F_{134} &= F_{143} = F_{145} = F_{154} = F_{215} = F_{251} = \frac{1}{2}, \\ F_{312} &= F_{314} = F_{321} = F_{323} = F_{332} = \frac{1}{2}, \\ F_{341} &= F_{415} = F_{451} = F_{525} = F_{552} = \frac{1}{2}, \\ F_{123} &= F_{125} = F_{132} = F_{152} = -\frac{1}{2}, \\ F_{345} &= F_{354} = F_{534} = F_{543} = -\frac{1}{2}, \\ F_{255} &= -F_{233} = 1.\end{aligned}$$

Using the general form of F , it can be seen that \mathfrak{g}_6 is of type $\mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_6 \oplus \mathcal{F}_9 \oplus \mathcal{F}_{10}$.

The nonzero components of Ricci curvature tensor are in the following:

$$\begin{aligned}Ric_{11} &= -1, Ric_{22} = -1, \\ Ric_{33} &= -\frac{1}{2}, Ric_{44} = \frac{1}{2}, \\ Ric_{55} &= \frac{1}{2}.\end{aligned}$$

With the aid of above relations, the scalar curvature tensor is $scal = -\frac{3}{2}$. The nonzero components of $\mathcal{L}_\xi g$ for the structure (3.2) are the following:

$$(\mathcal{L}_\xi g)_{23} = (\mathcal{L}_\xi g)_{32} = (\mathcal{L}_\xi g)_{34} = (\mathcal{L}_\xi g)_{43} = -1.$$

It is not hard to check that \mathfrak{g}_5 is neither Einstein-like nor Ricci-like soliton for structures (3.1) and (3.2).

Note that all components $\mathcal{L}_\xi g$ for the structure (3.1) are zero.

As a result, we get the followings.

Corollary 3.8. *The vector field ξ defined on the Lie algebras \mathfrak{g}_i for the structure (3.1) for $i = 1, \dots, 6$ is a Killing vector field.*

Corollary 3.9. *The structures given in (3.1) and (3.2) on five dimensional nilpotent Lie algebras \mathfrak{g}_i are not para-Sasaki-like.*

4. Conclusion

In this manuscript, we give two different Riemannian Π -structure on 5-dimensional nilpotent Lie algebras. The classes of given structures on 5-dimensional nilpotent Lie algebras are determined. We obtain the examples from certain classes. Only the \mathfrak{g}_1 Lie algebra among 5-dimensional nilpotent Lie algebras is Einstein-like and admit Ricci-like soliton according to the structure given in (3.1). 5-dimensional nilpotent Lie algebras \mathfrak{g}_i for the structure (3.2) are neither Einstein-like nor Ricci-like soliton.

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